

## Current Mirrors

### Basic BJT Current Mirror

Current mirrors are basic building blocks of analog design. Fig. 1 shows the basic npn current mirror. For its analysis, we assume identical transistors and neglect the Early effect, i.e. we assume  $V_A \rightarrow \infty$ . This makes the saturation current  $I_S$  and current gain  $\beta$  independent of the collector-base voltage  $V_{CE}$ . The input current to the mirror is labeled  $I_{REF}$ . This current might come from a resistor connected to the positive rail or a current source realized with a transistor or another current mirror. The emitters of the two transistors are shown connected to ground. These can be connected to a dc voltage, e.g. the negative supply rail.

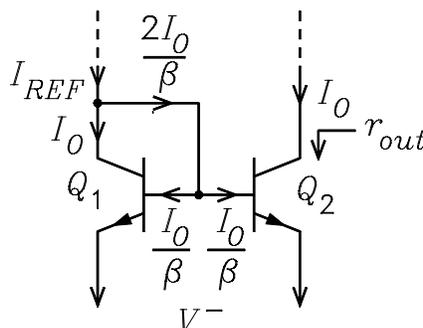


Figure 1: Basic current mirror.

The simplest way to solve for the output current is to sum the currents at the node where  $I_{REF}$  enters the mirror. Because the two transistors have their base-emitter junctions in parallel, it follows that both must have the same currents. Thus, we can write the equation

$$I_{REF} = I_O + \frac{2I_O}{\beta}$$

Solution for  $I_O$  yields

$$I_O = \frac{I_{REF}}{1 + 2/\beta}$$

Note that this equation predicts that  $I_O < I_{REF}$  unless  $\beta \rightarrow \infty$ . Because the Early effect has been neglected in solving for  $I_O$ , the output resistance is infinite. If we include the Early effect and assume that it has negligible effect in the solution for  $I_O$ , the output resistance is given by

$$r_{out} = r_{o2} = \frac{V_A + V_{CE2}}{I_O}$$

For a more accurate analysis, we can include the Early effect in calculating the output current. Consider the circuit in Fig. 2. If the transistors have the same parameters, we can write

$$\begin{aligned} I_{C1} &= I_{S0} \exp\left(\frac{V_{BE}}{V_T}\right) & I_{B1} &= \frac{I_{C1}}{\beta_0} \\ I_O &= I_{S0} \left(1 + \frac{V_{CB2}}{V_A}\right) \exp\left(\frac{V_{BE}}{V_T}\right) & I_{B2} &= \frac{I_O}{\beta_0 (1 + V_{CB2}/V_A)} \end{aligned}$$

By taking the ratio of  $I_O$  to  $I_{C1}$ , we obtain

$$I_O = \left(1 + \frac{V_{CB2}}{V_A}\right) I_{C1}$$

Note that the voltage  $V_{CB}$  is used for the Early effect, not the voltage  $V_{CE}$ . This agrees with the Ebers-Moll model for the BJT that is used in SPICE. In deriving the small-signal  $\pi$  and T models, the derivations are simplified by the use of  $V_{CE}$ . Thus  $V_{CE}$  is used in most contemporary texts for the Early effect. The correct value is  $V_{CB}$  which is used here.

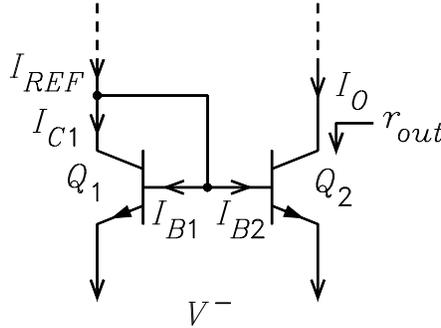


Figure 2: Circuit for including the Early effect.

Summing currents at the node where  $I_{REF}$  enters the circuit yields

$$I_{REF} = I_{C1} + \frac{I_{C1}}{\beta_0} + \frac{I_O}{\beta_0(1 + V_{CB2}/V_A)} = I_{C1} + \frac{2I_{C1}}{\beta_0}$$

Thus  $I_{C1}$  is given by

$$I_{C1} = \frac{I_{REF}}{1 + 2/\beta_0}$$

It follows that  $I_O$  is given by

$$I_O = \left(1 + \frac{V_{CB2}}{V_A}\right) I_{C1} = \frac{(1 + V_{CB2}/V_A)}{1 + 2/\beta_0} I_{REF}$$

Note that this equation predicts that  $I_O$  can be larger than  $I_{REF}$ . The output resistance is given above.

Note that the effect of a finite  $\beta$  is to reduce  $I_O$  but the effect of the Early effect is to increase it. Because of the Early effect, the output current can be greater than the input current. One way of obtaining a better match between the input and output currents is to use series emitter resistors on the transistors. If the current in one transistor increases, it causes the voltage across its emitter resistor to increase, which causes a decrease in its base-emitter voltage. This causes the current to decrease, thus causing the two transistors to have more equal currents. A typical value for the emitter resistors might be  $100\Omega$ . With these resistors,  $R_{te2}$  is no longer zero so that the output resistance is increased. It is given by  $r_{out} = r_{ic2}$  which can be much greater than  $r_{o2}$ .

### BJT Mirror with Base Current Compensation

Figure 3 shows the basic current mirror with a third transistor added. The collector of  $Q_3$  must be connected to a positive reference voltage, e.g. the positive supply rail, which biases it in the active

mode. If we neglect the Early effect and assume all transistors are identical, we can write

$$I_{REF} = I_O + \frac{2I_O}{\beta(1+\beta)}$$

Solution for  $I_O$  yields

$$I_O = \frac{I_{REF}}{1 + 2/[\beta(1+\beta)]}$$

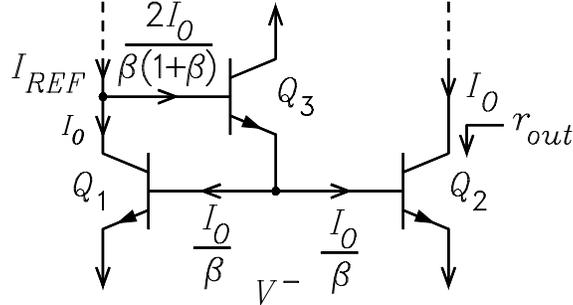


Figure 3: Mirror with base current compensation.

For a non-infinite Early voltage and  $V_{CB1} = V_{BE3} \ll V_A$ , it can be shown that the output current is given by

$$I_O = \frac{I_{REF} (1 + V_{CE2}/V_A)}{1 + 2/[\beta_0(1 + \beta_3)]}$$

where

$$\beta_3 = \beta_0 \left( 1 + \frac{V_{CE3}}{V_A} \right)$$

### BJT Wilson Mirror

A Wilson current mirror is shown in Fig. 4. We neglect the Early effect in the analysis and assume the transistors to have identical parameters. The emitter current in  $Q_3$  is  $I_O/\alpha$ . This current is the input to a basic current mirror consisting of  $Q_1$  and  $Q_2$ . This current is mirrored into the collector of  $Q_1$  by dividing by  $(1 + 2/\beta)$ . At the node where  $I_{REF}$  enters the mirror, we can write

$$I_{REF} = \frac{I_O/\alpha}{1 + 2/\beta} + \frac{I_O}{\beta} = I_O \frac{1 + \beta}{2 + \beta} + \frac{I_O}{\beta}$$

Solution for  $I_O$  yields

$$I_O = \frac{I_{REF}}{(1 + \beta)/(2 + \beta) + 1/\beta}$$

The advantage of the Wilson mirror over the current mirrors examined above is that it has a much higher output resistance. This is caused by two positive feedback effects. To see how this occurs, suppose a test current source is connected between the mirror output and ground. If the source delivers current to the output node, the voltage increases. This causes a current to flow through  $r_{o3}$ , causing the emitter voltage of  $Q_3$  and the base voltage of  $Q_1$  to increase. The increase in voltage at the emitter of  $Q_3$  causes its collector voltage to increase because  $Q_3$  is a common-base stage for an emitter input. Because  $Q_1$  is a common-emitter stage for a base input, the increase

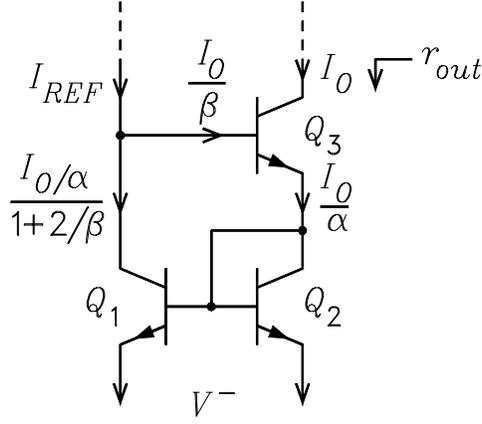


Figure 4: Wilson mirror.

in voltage at its base causes the collector voltage of  $Q_1$  and the base voltage of  $Q_3$  to decrease. Because  $Q_3$  is a common-emitter stage for a base input, the decrease in voltage at its base causes its collector voltage to increase. Thus there are two positive feedback effects which cause the collector voltage of  $Q_3$  to increase to a larger value. Because  $r_{out}$  is the ratio of the collector voltage of  $Q_3$  to the current in the test source, it follows that the output resistance is increased.

### BJT Low-Level Mirror

The circuit shown in Fig. 5 is a low-level current mirror. It can be used when it is desired to have a much lower output current than input current. For the analysis, we neglect the Early effect, assume identical transistors, and assume that  $\beta \rightarrow \infty$ . We can write

$$I_{REF} = I_S \exp\left(\frac{V_{BE1}}{V_T}\right) \quad I_O = I_S \exp\left(\frac{V_{BE1} - I_O R_E}{V_T}\right)$$

By taking ratios, we obtain

$$\frac{I_{REF}}{I_O} = \exp\left(\frac{I_O R_E}{V_T}\right)$$

This equation cannot be solved for  $I_O$ . If  $I_{REF}$  and  $I_O$  are specified, it can be solved for  $R_E$  to obtain

$$R_E = \frac{V_T}{I_O} \ln\left(\frac{I_{REF}}{I_O}\right)$$

As an example, suppose  $I_{REF} = 1 \text{ mA}$ ,  $V_T = 25 \text{ mV}$ , and  $I_O = 50 \mu\text{A}$ . It follows from this equation that  $R_E = 1498 \Omega$ . The effect of this large a value of  $R_E$  on  $r_{out}$  is to make it greater than  $r_0$ . To calculate  $r_{out}$ , we must know the small-signal Thévenin resistance  $R_{tb2}$  looking out of the base of  $Q_2$ . Note that  $Q_1$  is a bjt connected as a diode and exhibits a small-signal resistance  $r_{01} \parallel [r_{x1}/(1 + \beta_1) + V_T/I_{E1}] \simeq V_T/I_{E1} = 25 \Omega$ . This is in parallel with the small-signal resistance looking up into the  $I_{REF}$  source. Thus an upper bound on  $R_{tb2}$  is  $25 \Omega$ . Let us assume  $r_{02} = 40 \text{ k}\Omega$ ,  $r_{x2} = 0$ ,  $\alpha_2 = 0.995$ , and  $\beta_2 = 199$ . It follows that  $r_{ie2} = R_{tb2}/(1 + \beta) + \alpha V_T/I_{C2} = 497.6 \Omega$ . Thus  $r_{out}$  is given by

$$r_{out} = r_{ic2} = \frac{40000 + 497.6 \parallel 1498}{1 - 0.995 \times 1498 / (497.6 + 1498)} = 159.5 \text{ k}\Omega$$

This is larger than  $r_{02}$  by a factor of almost 4.

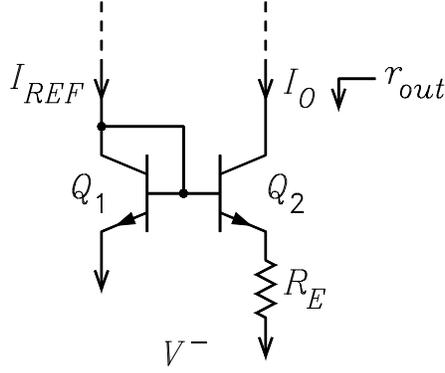


Figure 5: Low-level mirror.

## BJT Transconductance Op Amp

An example application of the current mirror is the transconductance op amp. The circuit is shown in Fig. 6. The circuit consists of an input diff amp and four Wilson current mirrors. For the analysis, we assume  $\beta \rightarrow \infty$  and  $V_A \rightarrow \infty$  for each bjt so that the output current from each mirror is equal to the input current. We assume that  $I_{ABC}$  splits equally between the emitters of  $Q_1$  and  $Q_2$ . Thus the total currents in  $Q_1$  and  $Q_2$ , respectively, are given by

$$i_{C1} = \frac{I_{ABC}}{2} + i_{c1} \quad i_{C2} = \frac{I_{ABC}}{2} + i_{c2} = \frac{I_{ABC}}{2} - i_{c1}$$

The latter expression for  $i_{C2}$  follows because  $i_{c1} + i_{c2} = 0$ .

It follows from the mirrored currents that the output current is given by

$$i_o = 2i_{c1}$$

If we neglect base currents and the Early effect,  $i_{c1} = i_{e1} = (v_{i1} - v_{i2}) / 2r_e$ , where  $r_e = 2V_T / I_{ABC}$ . Thus  $i_o$  is given by

$$i_o = \frac{I_{ABC}}{2V_T} (v_{i1} - v_{i2})$$

It can be seen that the transconductance gain is set by the current  $I_{ABC}$ . The gain can be varied by varying  $I_{ABC}$ . Because  $I_{ABC} \geq 0$ , the circuit operates as a two-quadrant multiplier. The circuit symbol for the transconductance op amp is shown in Fig. 7.

An example application of the transconductance op amp is a circuit which generates an amplitude modulated signal. The circuit is shown in Fig. 8. Let  $v_i$  and  $I_{ABC}$  be given by

$$v_i = V_1 \sin \omega_c t \quad I_{ABC} = I_Q (1 + m \sin \omega_m t)$$

where  $\omega_c$  is the carrier frequency,  $\omega_m$  is the modulating frequency, and  $m$  is the modulation index which must satisfy  $-1 < m < 1$ .

The current  $i_o$  is given by

$$i_o = \frac{I_{ABC}}{2V_T} \frac{v_i R_2}{R_1 + R_2}$$

If we assume that  $C_F$  is an open circuit at the operating frequencies, the current  $i_o$  must flow through  $R_F$ . Because the second op amp forces the voltage at its inverting input to be zero, the output voltage is given by

$$v_o = i_o R_F = \frac{I_{ABC}}{2V_T} \frac{v_i R_2}{R_1 + R_2} R_F = \frac{I_Q R_F}{2V_T} \frac{V_1 R_2}{R_1 + R_2} \sin \omega_c t (1 + m \sin \omega_m t)$$

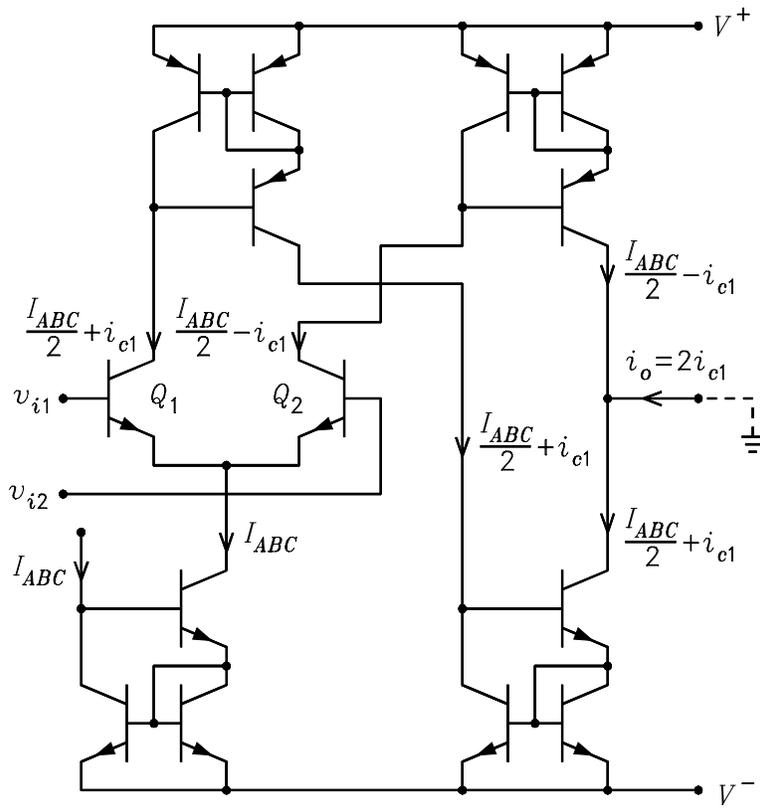


Figure 6: Transconductance op amp.

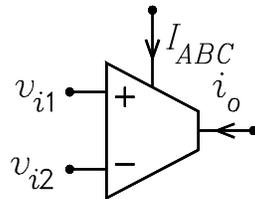


Figure 7: Transconductance op-amp symbol.

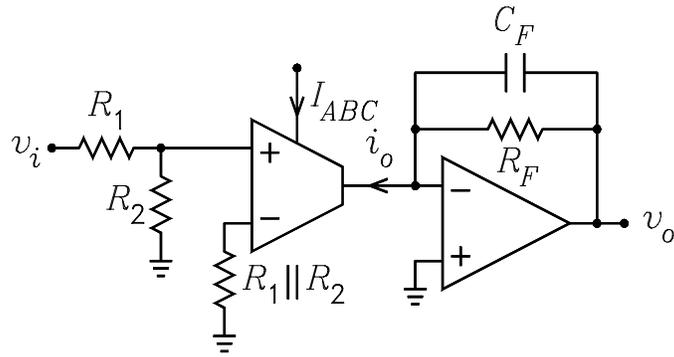


Figure 8: AM modulator.

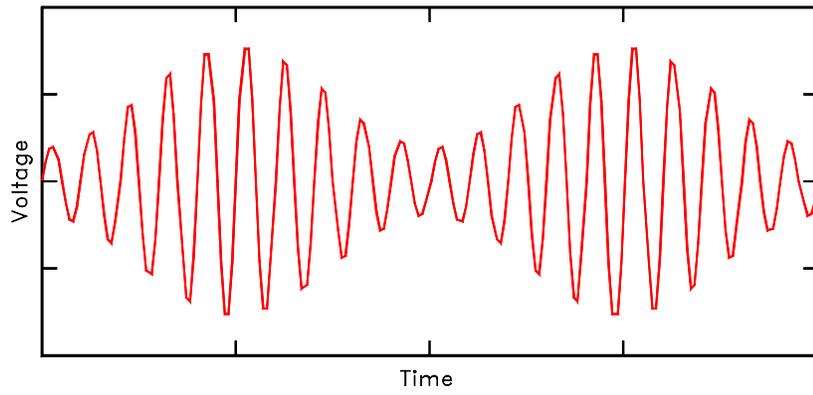


Figure 9: AM modulated waveform.

An example waveform for  $v_o$  for the case  $m = 0.6$  is shown in Fig. 9

The purpose of the voltage divider formed by  $R_1$  and  $R_2$  at the input to the transconductance op amp is to attenuate the input signal so that it does not overload the input differential amplifier. Ideally, the peak voltage should not exceed about 50 mV to prevent distortion in the output waveform. The resistor  $R_1 \parallel R_2$  is in series with the noninverting input so that both inputs to the differential amplifier see the same source resistance. Typically,  $R_1$  and  $R_2$  are chosen so that  $R_1 \parallel R_2$  is no larger than about  $100 \Omega$ . The capacitor  $C_F$  is necessary for proper high-frequency response. The capacitor must be chosen experimentally to prevent the gain of the circuit from peaking up at some high frequency where the circuit can oscillate. A method of determining the optimum value of  $C_F$  is to drive the circuit with a square wave for  $v_i$  and  $m = 0$  so that  $I_{ABC} = I_Q$ , i.e. a dc current. The capacitor can be experimentally adjusted to minimize any ringing on the output waveform without degrading the bandwidth of the circuit.